Weight-adjusted Bernstein-Bezier DG method for wave propagation in heterogeneous media

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North American High Order Methods Conference June 2-5, 2019 • Acoustic wave equation:

$$\frac{1}{c^2}\frac{\partial p}{\partial t} = \nabla \cdot \boldsymbol{u}, \qquad \frac{\partial \boldsymbol{u}}{\partial t} = \nabla p$$

• Elastic wave equation:

$$\rho \frac{\partial \boldsymbol{v}}{\partial t} = \sum_{i=1}^{d} \boldsymbol{A}_{i}^{T} \frac{\partial \boldsymbol{\sigma}}{\partial \boldsymbol{x}_{i}}, \qquad \boldsymbol{C}^{-1} \frac{\partial \boldsymbol{\sigma}}{\partial t} = \sum_{i=1}^{d} \boldsymbol{A}_{i} \frac{\partial \boldsymbol{v}}{\partial \boldsymbol{x}_{i}}$$

• Numerical scheme: high order, explicit time-stepping, parallelizable

- Unstructured (tetrahedral) meshes for geometric flexibility.
- High order: low numerical dissipation and dispersion.
- High order approximations: more accurate per unknown.
- Explicit time stepping: high performance on many-core.



Figure courtesy of Axel Modave.

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Graphics processing units (GPU).

Time-domain nodal DG methods

Assume $u(\mathbf{x}, t) = \sum \mathbf{u}_j \phi_j(\mathbf{x})$ on D^k

- Compute numerical flux at face nodes (non-local).
- Compute RHS of (local) ODE.
- Evolve (local) solution using explicit time integration (RK, AB, etc).



$$\frac{\mathrm{d}\mathbf{u}}{\mathrm{d}t} = \mathbf{D}_{\mathsf{X}}\mathbf{u} + \sum_{\mathsf{faces}} \mathbf{L}_{f}(\mathsf{flux}).$$

$$egin{aligned} m{M}_{ij} &= \int_{D^k} \phi_j(m{x}) \phi_i(m{x}) \ m{L}_f &= m{M}^{-1} m{M}_f. \end{aligned}$$

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- Bernstein-Bézier DG (BBDG): piecewise constant wavespeed
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High order approximation of media and geometry



• Piecewise constant wavespeed c^2 : efficient, but spurious reflections.

$$\frac{1}{c^{2}(\boldsymbol{x})}\frac{\partial \boldsymbol{p}}{\partial t} + \nabla \cdot \boldsymbol{u} = 0, \qquad \frac{\partial \boldsymbol{u}}{\partial t} + \nabla \boldsymbol{p} = 0$$

High order wavespeeds: weighted mass matrices. Stable, but requires pre-computation/storage of inverses or factorizations!

$$\mathbf{M}_{1/c^2} rac{\mathrm{d} \mathbf{p}}{\mathrm{d} t} = \mathbf{A}_h \mathbf{U}, \qquad \left(\mathbf{M}_{1/c^2}\right)_{ij} = \int_{D^k} rac{1}{c^2(\mathbf{x})} \phi_j(\mathbf{x}) \phi_i(\mathbf{x}).$$

• Weight-adjusted DG (WADG): energy stable approximation of M_{1/c^2}^k

$$\boldsymbol{M}_{1/c^2}^k rac{d\boldsymbol{p}}{dt} pprox \boldsymbol{M}^k \left(\boldsymbol{M}_{c^2}^k
ight)^{-1} \boldsymbol{M}^k rac{d\boldsymbol{p}}{dt} = \boldsymbol{A}_h \boldsymbol{U}$$

Reuses implementation for piecewise constant wavespeed

$$\frac{d\boldsymbol{p}}{dt} = \underbrace{\left(\boldsymbol{M}^{k}\right)^{-1}\left(\boldsymbol{M}_{c^{2}}^{k}\right)}_{\text{modified update}} \quad \underbrace{\left(\boldsymbol{M}^{k}\right)^{-1}\boldsymbol{A}_{h}\boldsymbol{U}}_{\text{RHS for } c=1}$$

• Modified update can be applied in a low storage manner using quadrature-based interpolation V_q and L^2 projection P_q .

Chan, Hewett, Warburton. 2017. Weight-adjusted discontinuous Galerkin methods: wave propagation in heterogeneous media.

Quadrature-based operators

• Using quadrature rule

$$\left(\boldsymbol{M}_{c^{2}}^{k}\right)_{ij} = \int_{D^{k}} c^{2}(\boldsymbol{x})\phi_{i}^{k}(\boldsymbol{x})\phi_{j}^{k}(\boldsymbol{x})d\boldsymbol{x} = J^{k}\sum_{n=1}^{N_{q}} c^{2}(\boldsymbol{x}_{n}^{q})\phi_{i}(\hat{\boldsymbol{x}}_{n}^{q})\phi_{j}(\hat{\boldsymbol{x}}_{n}^{q})\boldsymbol{w}_{n}$$

where $\mathbf{x}^q, \hat{\mathbf{x}}^q$ denote quadrature points on D^k and \hat{D} , respectively.

Writing *M*^k_{c²} into matrix form

$$oldsymbol{M}_{c^2}^k = J^k oldsymbol{V}_q^T$$
diag $(oldsymbol{w})$ diag $(c^2) oldsymbol{V}_q$

where

$$(\boldsymbol{V}_q)_{ij} = \phi_j(\hat{\boldsymbol{x}}_i^q)$$

• The modified update can be written as

$$\left(\boldsymbol{M}^{k}\right)^{-1}\left(\boldsymbol{M}_{c^{2}}^{k}\right) = \underbrace{\boldsymbol{M}^{-1}\boldsymbol{V}_{q}^{T}\operatorname{diag}\left(\boldsymbol{w}\right)}_{\boldsymbol{P}_{q}}\operatorname{diag}\left(c^{2}\right)\boldsymbol{V}_{q}$$

- V_q : evaluates function values at quadrature points.
- P_q : projects a function onto a polynomial space in L^2 sense.

Wave simulations in heterogeneous media



- L^2 convergence between optimal $O(h^{N+1})$, provable $O(h^{N+1/2})$.
- Difference between standard DG and WADG is $O(h^{N+2})$.

Chan, Hewett, Warburton. 2017. Weight-adjusted DG methods: wave propagation in heterogeneous media.

$$\frac{d\boldsymbol{p}}{dt} = \underbrace{\boldsymbol{P}_{\boldsymbol{q}} \text{ diag}(c^2) \boldsymbol{V}_{\boldsymbol{q}}}_{\text{modified update}} \underbrace{\left(\boldsymbol{M}^k\right)^{-1} \boldsymbol{A}_h \boldsymbol{U}}_{\text{RHS for } c=1}$$

- RHS for c = 1 produces a polynomial u.
- V_q evaluates values of u at quadrature points.
- Applying diag (c^2) to u gives the product c^2u .
- P_q projects $c^2 u$ onto a polynomial space of degree N.

$$\frac{d\boldsymbol{p}}{dt} = \underbrace{\boldsymbol{P}_{q} \text{ diag } (c^{2}) \boldsymbol{V}_{q}}_{\text{modified update}} \underbrace{\left(\boldsymbol{M}^{k}\right)^{-1} \boldsymbol{A}_{h} \boldsymbol{U}}_{\text{RHS for } c=1}$$

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Bernstein polynomial bases on simplices



Each function attains its maximum at an equispaced lattice point of a *d*-simplex.

- Bernstein polynomials are associated with vertex, edge, face, and interior equispaced nodes.
- Jumps across interfaces can be computed similarly to nodal bases (reuse nodal DG framework).

Bernstein polynomial

• On the reference tetrahedron, the barycentric coordinates are $\lambda_0 = -\frac{(1+r+s+t)}{2}, \ \lambda_1 = \frac{(1+r)}{2}, \ \lambda_2 = \frac{(1+s)}{2}, \ \lambda_3 = \frac{(1+t)}{2}.$

• The Nth degree Bernstein basis is defined as

$$B_{\alpha}^{N} = C_{\alpha}^{N} \lambda_{0}^{\alpha_{0}} \lambda_{1}^{\alpha_{1}} \lambda_{2}^{\alpha_{2}} \lambda_{3}^{\alpha_{3}}, \quad C_{\alpha}^{N} = \frac{N!}{\alpha_{0}! \alpha_{1}! \alpha_{2}! \alpha_{3}!},$$

where $\alpha = (\alpha_{0}, \alpha_{1}, \alpha_{2}, \alpha_{3})$ and $\alpha_{0} + \alpha_{1} + \alpha_{2} + \alpha_{3} = N.$

Simple degree elevation of Bernstein polynomials

$$B_{\alpha}^{N-1} = \sum_{j=0}^{d} \frac{\alpha_j + 1}{N} B_{\alpha+e_j}^{N}$$

leads to sparse one-degree elevation operators.

- Nodal DG: $O(N^6)$ cost in 3D vs $O(N^3)$ degrees of freedom.
- Switch to Bernstein basis: sparse and structured matrices.
- Optimal O(N³) application of differentiation and lifting matrices.



Nodal bases in one, two, and three dimensions.

Chan, Warburton 2015. GPU-accelerated Bernstein-Bezier discontinuous Galerkin methods for wave propagation (SISC).

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Sparse Bernstein differentiation matrices for the reference tetrahedron.

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Optimal $O(N^3)$ complexity "slice-by-slice" application of Bernstein lift.

Chan, Warburton 2015. GPU-accelerated Bernstein-Bezier discontinuous Galerkin methods for wave propagation (SISC).

Numerical experiments are implemented on a mesh with 61276 elements



Chan, Warburton. 2017. GPU-accelerated Bernstein-Bézier discontinuous Galerkin methods for wave problems.

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BBWADG: polynomial multiplication and projection



- BBWADG reuses the volume and surface kernels from BBDG
- Instead of using quadrature-based L^2 projection, BBWADG approximates $c^2(\mathbf{x})$ with degree M polynomial, use fast Bernstein algorithms for polynomial multiplication and projection.

Bernstein polynomial multiplication

• Assume $h(\mathbf{x}) = f(\mathbf{x})g(\mathbf{x})$, where f, g are Bernstein polynomials

$$f=c_1\lambda_0+c_2\lambda_1+c_3\lambda_2+c_4\lambda_3,\qquad g=\sum_{|mlpha|\leq N}d_{mlpha}B_{mlpha}^N$$

• We can split the multiplication into *d* + 1 parts. For example, the first part is

$$c_1 \sum_{|\alpha| \le N} d_{\alpha} \left(B_{\alpha}^N \lambda_0 \right) = c_1 \sum_{|\alpha| \le N} d_{\alpha} \frac{\alpha_j + 1}{N + 1} B_{\alpha + \mathbf{e}_0}^{N+1}$$

Bernstein polynomial multiplication



Bernstein polynomial multiplication: for fixed M, $O(N^3)$ complexity.

Fast Bernstein polynomial projection

- Given c²(x)u(x) as a degree (N + M) polynomial, apply L² projection matrix P^{N+M}_N to reduce to degree N.
- Polynomial L^2 projection matrix P_N^{N+M} under Bernstein basis:

$$\boldsymbol{P}_{N}^{N+M} = \underbrace{\sum_{j=0}^{N} c_{j} \boldsymbol{E}_{N-j}^{N} \left(\boldsymbol{E}_{N-j}^{N} \right)^{T}}_{\boldsymbol{\tilde{P}}_{N}} \left(\boldsymbol{E}_{N}^{N+M} \right)^{T}$$

• "Telescoping" form of \tilde{P}_N : $O(N^4)$ complexity, more GPU-friendly.

$$\left(c_0 \boldsymbol{I} + \boldsymbol{E}_{N-1}^{N}\left(c_1 \boldsymbol{I} + \boldsymbol{E}_{N-2}^{N-1}\left(c_2 \boldsymbol{I} + \cdots\right)\left(\boldsymbol{E}_{N-2}^{N-1}\right)^{T}\right)\left(\boldsymbol{E}_{N-1}^{N}\right)^{T}\right)$$

Illustration of GPU algorithm for \ddot{P}_N



BBWADG: approximating c^2 and accuracy



Approximating smooth $c^2(\mathbf{x})$ using L^2 projection: $O(h^2)$ for M = 0, $O(h^4)$ for M = 1, $O(h^{M+3})$ for $0 < M \le N - 2$.

BBWADG vs WADG (acoustic)



BBWADG vs WADG (elastic)



	N = 3	<i>N</i> = 4	N = 5	<i>N</i> = 6	<i>N</i> = 7	N = 8
WADG	2.02e-8	4.91e-8	1.20e-7	2.19e-7	4.87e-7	5.25e-6
BBWADG-1	2.09e-8	3.32e-8	6.56e-8	8.54e-8	1.35e-7	1.65e-7
Speedup	0.9665	1.4789	1.8292	2.5644	3.6074	31.8182

For $N \ge 8$, WADG becomes much more expensive because of the quadrature points construction.



Summary and acknowledgements

- Weight-adjusted DG: stability and efficiency for heterogeneous media.
- BBWADG: improved complexity for approximate wavespeeds.
- This work is supported by the National Science Foundation under DMS-1712639 and DMS-1719818.

Thank you! Questions?



Guo, Chan. 2018. Bernstein-Bezier weight-adjusted discontinuous Galerkin methods for wave propagation. Chan, Hewett, Warburton. 2016. Weight-adjusted DG methods: wave propagation in heterogeneous media (SISC). Chan 2017. Weight-adjusted DG methods: matrix-valued weights and elastic wave prop. in heterogeneous media (IJNME). Chan, Warburton 2015. GPU-accelerated Bernstein-Bezier discontinuous Galerkin methods for wave propagation (SISC).